

Continuous-Flow Synthesis of Orange Emitting Sn(II)-Doped CsBr Materials

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An ongoing demand toward lead-free all-inorganic cesium metal halide perovskites has presented Sn(II) as an ideal substitute of Pb(II) for applications in optoelectronic devices. The major concern regarding Sn(II) is the instability due to the ambient oxidation to Sn(IV). To expand the scope of traditional perovskite and analogues, herein the synthesis and optical performance of Sn(II)-doped CsBr, a new material formed by interstitial doping of Sn(II) into the CsBr matrix, are reported for the first time. This material is prepared following an antisolvent mediated recrystallization method using a continuous flow reactor, which is beneficial for scaling up the production compared to traditional batch reactors. Sn(II)-doped CsBr exhibits broadband orange emission with full-width-half-maximum of 180 nm and a photoluminescence quantum yield of 21.5%. The emission turned to be highly stable over 7 months despite containing Sn(II). It is suggested that this is due to interstitial location of Sn(II) atoms in bulk of microcrystals. A broadband emission and high aerobic stability are attractive properties of the material for white-light emitting applications.

1. Introduction

In the past decade, lead halide perovskites are on the edge-of-the-seat for their exciting power conversion efficiencies and

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display technologies,^[1–6] but their commercialization still faces significant challenges, as the toxicity of carcinogenic lead metal.^[7,8] As a suitable alternative, ongoing research is carrying out on lead-free metals such as Sn, Bi, Ag, Cu, In, and Sb-based perovskites and double perovskites, which also has proven their exceptional optoelectronic properties and promising solar to power energy conversion efficiencies.^[9–18] Among the different halide perovskites the fully inorganic are receiving an increasing attention due to their higher stability.^[19,20] All-inorganic cesium tin halides present a variety of compositions, viz. CsSnX₃, Cs₄SnX₆, Cs₂SnX₆ have been explored with good optoelectronic properties and photovoltaics.^[21–23] Importantly, Sn(II) is isovalent to Pb(II) and CsSnX₃/Cs₄SnX₆ is isostructural with CsPbX₃/Cs₄PbX₆. Hence tin can

be considered as a satisfactory replacement to explore lead-free metal halide perovskites and analogues.^[24] Unfortunately, there is a limitation in tin-based perovskites owing to the oxidation of Sn²⁺ to Sn⁴⁺ upon exposure to ambient air and moisture, which degrades the material by creating trap states that irreversibly collapse the optical emission (and/or carrier lifetimes) of the material. This phenomenon somehow limits the commercial applicability of tin-based perovskites. To overcome this issue, the replacement of Sn²⁺ by Sn⁴⁺ was proposed as an alternative to get Cs₂SnX₆ with high valency of tin (Sn⁴⁺). This typical structure of perovskites is referred to as vacancy-ordered double perovskites.^[25]

In light of the emissive nature of the corresponding tin halide perovskites, CsSnX₃ exhibits an excitonic emission with narrow full-width-half-maximum (FWHM). Unfortunately, they possess a very low photoluminescence quantum yield (PLQY) in comparison to the lead halide perovskites.^[26] However, current sign of progress on 0D tin halide perovskites have jumped up to importance in the field of optoelectronics due to their satisfactory photoluminescence with broad coverage across the visible region.^[23] Recently, several reports on 0D Cs₄SnBr₆ published portraying a self-trapped excitonic (STE) emission originated from the isolated [SnBr₆]⁴⁻ octahedra.^[27–31] Typical characteristics of STE emissions are i) broad emission and ii) large stokes shift.^[32] Hence, the generation of white light due to covering a broad optical window and the reabsorption-related

optical losses are ruled out here. However, these materials have some limitations in terms of broad application and commercialization because of their ambient instability. These materials transform to Cs_2SnBr_6 in presence of air, subsequently lose the emission.^[28] Cs_2SnX_6 , vacancy ordered double perovskites are a class of materials, where the Sn atom has a +4-oxidation state and hence they possess higher stability in the open air.^[18] Their optical performances enable many application segments including photovoltaics and light-emitting materials. Unfortunately, the potential optoelectronic application is also break-off here due to experiencing very low PLQY.^[18,33]

For the preparation of cesium tin halide perovskite and analogues, antisolvent mediated recrystallization is the easiest method followed previously.^[29] Nevertheless, the labor intensive batch methodology is cumbersome, time consuming and presents limited scalability. Hence, we choose flow-reactors to have an accurate control over the crystallization conditions by finely tuning the flow rates of the precursors and antisolvent solutions. Flow-reactors have been used to synthesize and to crystallize nanomaterials and APbX_3 (A = formadimium/methylammonium/Cs) perovskite nanocrystals and it was established that these methods are favorable for higher reproducibility and scalable production.^[34–40] Continuous flow system can be modulated controlling flow rates, temperatures and pressures to mix a homogeneous precursor solution/organic ligand(s)/antisolvent in micromixers and microreactors, with the possibility to finely tune residence times. Herein, we have prepared Cs_4SnBr_6 and Sn(II)-doped CsBr following a simple flow-reactor by setting down commands of different precursor and antisolvent flow rates. The former material is well-studied in accordance to their STE emission, A-site doping, and thermographic applications.^[28–31,41] We found Cs_4SnBr_6 is unstable and it degrades in open air over time and eventually transforms to Cs_2SnBr_6 . However, the other material, Sn(II)-doped

CsBr was prepared for the first time, which exhibits broad-range STE emission (PLQY = 21.5%, FWHM \approx 180 nm) and highly stable in the open air over 7 months as per laboratory report.

2. Result and Discussions

For the material synthesis the room-temperature antisolvent mediated recrystallization method was followed, where dimethyl formaldehyde (DMF) was used to dissolve the precursors and toluene to precipitate the crystalline material. The fundamental of this process includes dissolving precursor solutes into a polar (or good) solvent and by the incorporation of an antisolvent (or poor) dissolved ions arrange among themselves to form an energetically favorable crystalline material. We modified the synthetic procedure by changing from batch to continuous flow. Two high performance liquid chromatography pumps were used to flow the precursors dissolved in DMF and the toluene to induce the crystallization. These two pumps were connected with 1/16" T-mixer (Swagelok), followed by a short tubular microreactor, which served to rapidly mix the precursor solution and the antisolvent, followed by the collection of the crystalline powder at the outlet. The schematic of the flow-reactor is displayed on Figure 1a. The flow rate of the antisolvent was fixed to 1 mL min^{-1} and the flow rate of the precursor solution was changed from 0.1 to 0.5 mL min^{-1} . Interestingly, we obtained two different products at different precursor flow rates (depicted in Figure S1, Supporting Information). We observed the preferential formation of one product at the lower flow rates of precursors solution ($0.1\text{--}0.2 \text{ mL min}^{-1}$) and another at higher flow rates ($0.3\text{--}0.5 \text{ mL min}^{-1}$), respectively. At lower flow rates, Cs_4SnBr_6 , a white-colored air-unstable green-emitting material was formed, whose emission was quenched

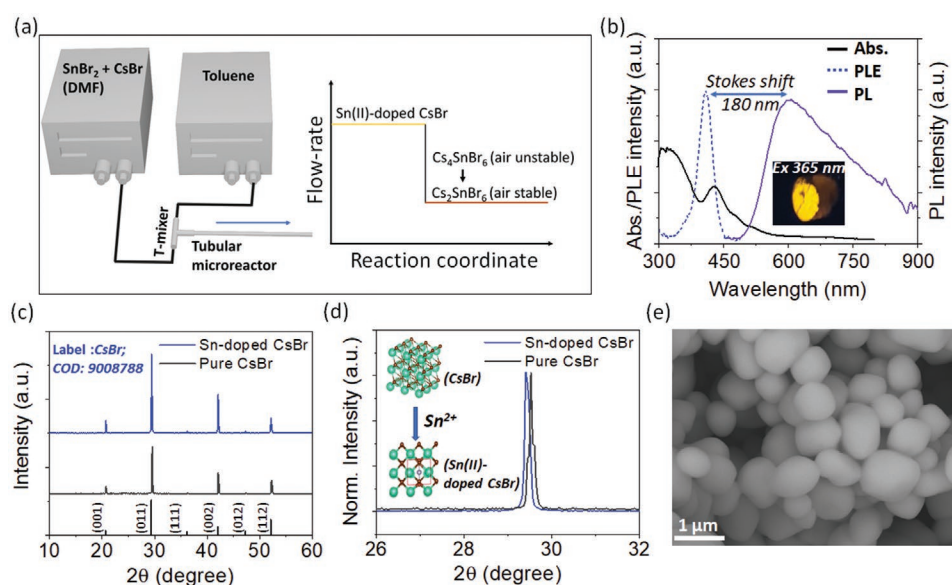


Figure 1. a) Schematic presentation of experimental setup and typical reaction procedure. b) UV–vis absorbance (Abs.), PL, and PLE spectra of Sn(II)-doped CsBr powder. Inset shows illuminated orange emission of Sn(II)-doped CsBr under 365 nm UV-lamp. c) XRD patterns of Sn(II)-doped CsBr and undoped CsBr. The label of XRD for matching is referenced from COD: 9008788. d) Enlarged view of (011) peak to show the peak shifting belongs to lattice expansion for Sn(II)-dopant incorporation into the CsBr lattice. e) SEM image of as-prepared Sn(II)-doped CsBr powder.

during the drying process and transformed to CsBr and Cs_2SnBr_6 in the open air. The white-colored Cs_4SnBr_6 turned to orange-colored due to the formation of Cs_2SnBr_6 in the open air. At higher flow rates, Sn(II)-doped CsBr, a yellowish-colored powder was precipitated, which possessed a strong orange emission found to be stable over 7 months. The digital images of Sn(II)-doped CsBr powder and the oxidized powder (Cs_4SnBr_6 , Cs_2SnBr_6) are shown in Figure S2a,b in the Supporting Information. The illuminated images of orange emitting Sn(II)-doped CsBr, Cs_4SnBr_6 (under inert atmosphere) and their poly(methyl methacrylate) (PMMA)-coated films on top of glass slides under the 365 nm UV-lamp are provided in Figure S3a–c in the Supporting Information.

Figure 1b presents the optical absorbance (Abs.), photoluminescence (PL), and PL excitation spectrum (PLE) of as-obtained Sn(II)-doped CsBr. The figure inset shows the digital image of the orange light emission from Sn(II)-doped CsBr powders under a 365 nm excitation. For PL measurements, the material was excited at 420 nm, and for PLE, the detection wavelength was fixed at 600 nm. This emission exhibits large FWHM (≈ 180 nm) and Stokes shift (180 nm).

The X-ray diffraction (XRD) pattern of Sn(II)-doped CsBr synthesized at high flow rate (0.4 mL min^{-1}) exactly matched with cubic CsBr ($Fm-3m$ space group). A control reaction was carried out without employing SnBr_2 into the reaction and as-obtained undoped material was characterized as the same XRD pattern (see Figure 1c). The crystalline phase is matched with CsBr COD:9008788. A XRD peak shifting toward the

lower angle was observed indicating a lattice expansion (see Figure 1d). From the theoretical calculations (explained in a later section), it is shown that doping Sn^{2+} is incorporated at the interstitial position of the CsBr lattice, in good agreement with the lattice expansion observed in Figure 1d. The morphology of the as-obtained Sn(II)-doped CsBr material was monitored by scanning electron microscope (SEM) and transmission electron microscopic (TEM) imaging (Figure 1e and Figure S4, Supporting Information), which shows quasi spherical particles ($300 \text{ nm}^{-1} \mu\text{m}$).

We have systematically studied the formation of materials under different precursor solution flow rates and we analyzed the results by UV–vis absorption spectroscopy (Figure 2a). The flow rate ratio is defined as X/Y where X and Y are the precursor ($\text{SnBr}_2 + \text{CsBr}$) and the antisolvent (toluene) flow rates, respectively. At lower flow rate ratios (0.1/1 and 0.2/1), there is an absorbance peak near 340 nm, which belongs to 0D Cs_4SnBr_6 perovskite, referenced from the literature report.^[27] For the sample prepared at 0.2/1 ratio, an evolution of the broad absorbance peak near 430 nm is observed, which is due to the formation of Sn(II)-doped CsBr, see Figure 2a. The 430 nm peak intensifies and the 340 nm peak disappears finally with increasing flow rate ratio that confirms the formation of only Sn-doped CsBr product at high flow rate ratios (green and blue lines Figure 2a for 0.3/1–0.4/1 and Figure S5a, Supporting Information, for 0.5/1). Corresponding PL and PLE spectra of these sets of materials are presented in Figure 2b. The materials were excited at 420 nm to detect PL. The materials

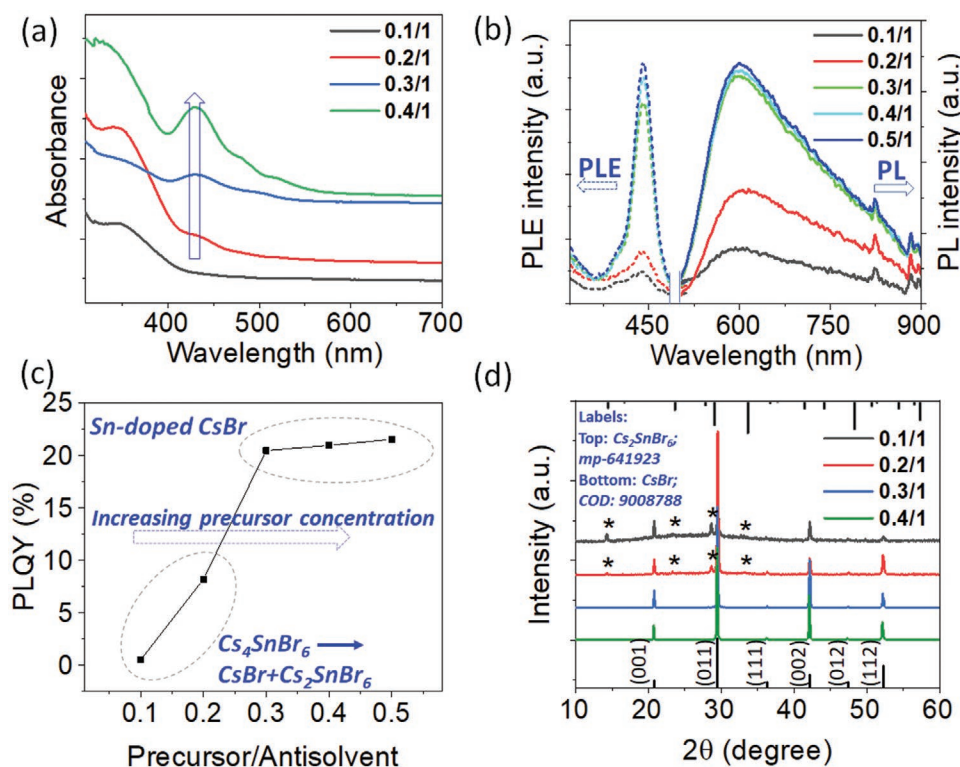


Figure 2. a) UV–vis absorbance and b) PL (solid lines) and PLE (dashed lines) of samples prepared at different flow rates. PL spectra were obtained under the excitation of 420 nm. c) A plot of PLQY versus precursor/antisolvent flow ratio. Samples were excited at 420 nm for the PLQY measurements. d) XRD patterns of powder samples prepared at different flow rates. Peaks belonging to Cs_2SnBr_6 are marked with asterisk.

prepared at low flow rate ratios exhibit a low intensity PL, that increases at higher flow rate ratios, stabilizing the intensity for flow rate ratios higher than 0.3/1. The orange-emitting material shows a significantly narrow PL excitation range between 390 and 450 nm, beyond which the material is barely luminescent or even nonluminescent at all, which is unusual for band-to-band emitting semiconductors. The PLQs of these samples were measured and plotted versus the flow rate ratio that is presented in Figure 2c, where a highest PLQY of 21.5% was observed at higher flow rate ratios. All the obtained PLQY values are provided in Table S1 in the Supporting Information. The excitation wavelength was 420 nm for the PLQY measurements. The poor orange PL obtained at lower flow rate ratios is due to the preferential crystallization of Cs_4SnBr_6 , which has a green emission under the excitation of 337 nm, see Figure S6 in the Supporting Information, and a minimum formation of Sn(II)-doped CsBr. Table S2 in the Supporting Information provides the corresponding PLQY of the green emission from Cs_4SnBr_6 (excitation wavelength 337 nm). Contrarily, it was found that the intensity of the green emission decreased with increasing precursor flow rate ratio (Figure S6a, Supporting Information).

The as-prepared samples were dried in the ambient conditions for the XRD characterization, presented in Figure 2d. The XRD patterns of all of these samples were matched with CsBr ($Pm-3m$ space group, cubic phase). At lower flow rates, additional peaks were observed (marked with asterisk in Figure 2d, first two patterns from the top), which belonged to the Cs_2SnBr_6 compound (mp-641923, triclinic phase). At higher flow rate ratios (0.3/1–0.5/1), no additional peaks were

observed for any of the samples corresponding to the phase-pure CsBr cubic lattice without any impurities (Figure 2d, last two patterns at the bottom for 0.3/1–0.4/1 and Figure S5b, Supporting Information, for 0.5/1). In an alternative way, the unstable green-emitting product obtained at 0.1/1 was collected under the N_2 and it was coated with PMMA to prevent it from ambient atmosphere. Figure S7 in the Supporting Information presents the XRD pattern of the PMMA-coated green emitting product, which matches with the Cs_4SnBr_6 crystal phase. The ambient instability of the green-emitting Cs_4SnBr_6 is due to the oxidation of Sn(II) to Sn(IV) and it transformed to Cs_2SnBr_6 and CsBr.^[28] The powder samples prepared at low flow rates only have the XRD patterns for Cs_2SnBr_6 (do not confuse with PMMA-coated sample). Beyond 0.3/1, the XRD patterns only belong to CsBr, confirms the formation of phase-pure Sn(II)-doped CsBr. Indeed, it was discarded the possibility of the new material being Cs_2SnBr_6 , the oxidized product of Cs_4SnBr_6 , by separately synthesizing Cs_2SnBr_6 by wet-chemical method (see the Experimental Section), and comparing the XRD patterns (Figure S8, Supporting Information). This confirmed that Cs_2SnBr_6 phase shows cubic phase and $Fm-3m$ space group (mp-641923) as a difference from the new material. This information explained the phase of orange emitting material is identical to CsBr and not belongs to Cs_2SnBr_6 .

Raman spectroscopic measurements were carried out to study the Sn–Br vibration into the crystal lattice. Herein, comparing the as-obtained spectra of the dried powder samples obtained at 0.1/1 and 0.5/1 are presented in Figures 3a,b, respectively. Surprisingly, we observed a discrepancy in the Raman spectrum pattern depending on the flow rate ratios.

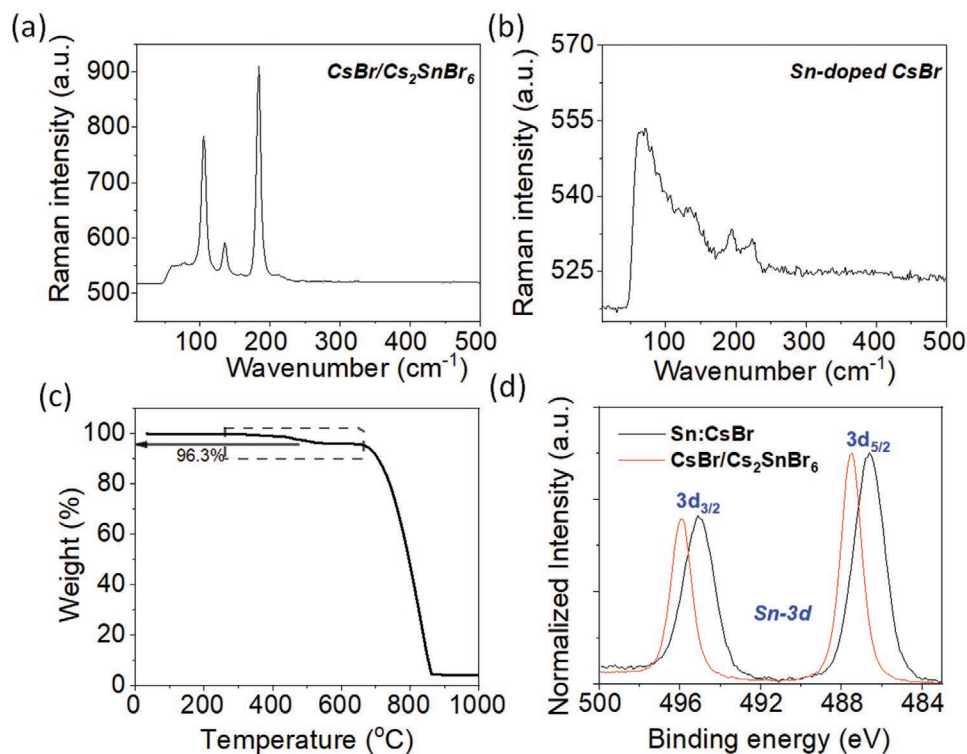


Figure 3. Raman spectra of a) $\text{CsBr}/\text{Cs}_2\text{SnBr}_6$ (oxidized from Cs_4SnBr_6) and b) Sn(II)-doped CsBr powder samples. c) TGA of Sn(II)-doped CsBr powder. d) XPS of Sn-3d of Sn(II)-doped CsBr (black) and $\text{CsBr}/\text{Cs}_2\text{SnBr}_6$ (red; oxidized from Cs_4SnBr_6).

At 0.1/1, there are three intense peaks at 106.9, 136.1, and 184.3 cm^{-1} while those are absent for the sample prepared at 0.5/1. As discussed earlier, due to ambient oxidation, green-emitting Cs_4SnBr_6 transformed to CsBr and Cs_2SnBr_6 . In Cs_2SnBr_6 , there are three Raman active vibrational modes that are present due to the vibrations inside the $[\text{SnBr}_6]^{2-}$ octahedra, which are the Br–Sn–Br asymmetric bending, Sn–Br asymmetric stretching, and symmetric Sn–Br stretching vibration.^[42] In the case of Sn-doped CsBr, Sn(II) is occupying interstitial sites, as we discuss below, in the CsBr crystal lattice and conjugated to the Br-atoms linearly (Figure S9, Supporting Information). Hence, such vibrational modes are absent for the interstitial SnBr_2 (doped) units into the CsBr lattice.

Thermogravimetric analysis (TGA) presented in Figure 3c also predicts a similar result for the presence of a minimum content of Sn, as Sn-doped CsBr has a 3.7% mass loss near 530 °C and rest is decomposed at 780 °C. A control experiment of undoped CsBr confirms the rest of the mass lost appertain to CsBr (see Figure S10, Supporting Information). The initial loss for the dopant ion inclusion now led us to calculate the dopant ion percentage into the as-prepared material and hence, we have calculated 10% of initial mass loss for the material synthesized at 0.1/1, which indeed Cs_4SnBr_6 transformed to CsBr and Cs_2SnBr_6 (see Figure S11, Supporting Information). Hence, from the synthesis point of view, a high amount of SnBr_2 was precipitated at low precursor flow rates, whereas the tin content decreased in the final material with the increase of precursor concentration (simultaneously, the DMF content). Elemental composition was further verified using energy dispersive X-ray spectroscopy (EDS) analysis from the TEM measurements, which estimates the Sn-percentage from 1.3–3.8% (Figure S12 and Table S3, Supporting Information).

X-ray photoelectron spectroscopy (XPS) was carried out to study the oxidation state of Sn-dopants into the CsBr crystal. Figure 3d presents the Sn-3d XPS peaks of Sn-doped CsBr (black) and $\text{CsBr}+\text{Cs}_2\text{SnBr}_6$ (red) oxidized from the initial formation of Cs_4SnBr_6 . Binding energy of Sn-3d_{5/2} for Sn-doped CsBr appears at 486.6 eV ascribed to Sn^{2+} in Sn-doped CsBr whereas $\text{CsBr}+\text{Cs}_2\text{SnBr}_6$ appears at higher binding energy 487.5 eV associated to Sn^{4+} , which was expected due to the oxidized powder.^[43]

Eventually, in order to confirm the exact nature of the material prepared at high flow ratios, the system was further investigated by density functional theory (DFT) calculations performed with Quantum ESPRESSO,^[44] a plane-wave periodic DFT code. First, we geometrically relaxed a clean unit cell composed of 54 CsBr atoms using a single k point (Γ for the direct bandgap), spin–orbit interactions, and PBEsol pseudopotentials generated with the generalized gradient approximation (GGA) for the exchange–correlation functional. The Cs–Br bond length in the geometrically relaxed atomic structure is 3.80 Å, 2.2% larger than the experimental 3.72 Å. Secondly, we doped the unit cell with an interstitial Sn atom placed between two Br atoms at the central region of a quasi-orthorhombic void space, see Figure 4a, and relaxed the atomic structure again, resulting in a Sn–Br distance of 2.94 Å. The considered 3.7% doping ratio was previously estimated in an experiment consisting in heating synthesized powder samples of Sn-doped CsBr and measuring the weight loss (see Figure 3c). The calculated densities of states (DOS) for this system, addressed in Figure 4a, reveal a bandgap

of 4 eV which underestimates the experimental 7.3 eV,^[45] as usual when DFT computations are run at the GGA level, however allows to observe the trends caused by Sn addition. The DOS show as well two peaks of p-like character within the bandgap, appended by the dopant: the left-hand one is formed by two polarons at two degenerate energy levels, indicating a +2 oxidation state for the Sn atom, and the right-hand one is empty and also doubly degenerate. The unoccupied peak is placed at 2.7 eV above the valence band minimum (VBM); this energy value is close and consistent with the 2.9 eV (430 nm) peak obtained in the experimental absorption spectrum. Furthermore, the occupied lowest-energy states show a characteristic ring shape (quantum numbers $l = +1$ and $m = -1$), and the empty states in the second peak are also ring shaped ($l = +1$ and $m = +1$). Since impurity states are strongly localized at the Sn site, the photoexcitation of one electron from the VBM into these states would result in self-trapped excitons, that is excitons hooked to the crystal lattice. Moreover, by looking at the DOS of Figure 4a and the local contributions of the Sn dopant and the two Br neighboring atoms, we conclude that the Sn atomic orbitals are only slightly hybridized with the host crystal states, in consonance with the interstitial nature of the dopant.

We have further obtained theoretical XRD patterns for the clean and doped cases. The deviation of the main peak in the XRD pattern of the doped structure toward the left relative to the same peak in the spectrum of the clean host crystal is 0.4°. This value is larger than the 0.12° shift obtained in the XRD experiments (Figure 1d), what would indicate that the real concentration of Sn dopants in the synthesized powder samples is smaller than the initially estimated 3.7%.

PL spectra of Sn(II)-doped CsBr measured in 200–330 K are shown in Figure 5a. Temperature increase in the indicated diapazon results in approximately twofold decrease of PL integral intensity, which is accompanied by similar twofold decrease of the main (shorter) component lifetime for kinetics measured at 600 and 700 nm (Figures S13 and S14, Supporting Information). The found temperature dependence can be explained by freezing nonradiative deactivation channels at lower temperatures, which are activated with temperature. Another interesting aspect regarding stability is presented in Figure 5b. The material exhibits long-term optical stability and hence an identical emission intensity even after the 7 months. This material belongs to the Sn(II) family and hence it has extended long-term stability and possibly in competition with lead halide perovskites. Presumably, the extended stability is due to i) doping at the interstitial position not facing the surface and ii) protective covering by the Cs–Br crystal matrix.

To further elucidate this newly explored orange emitting material for light-emitting diode (LED) applications, herein, we dip-coated Sn-doped CsBr on top of a commercially available 430 nm blue chip. The electroluminescence spectra were detected at different voltage and presented in Figure 5c. The blue electroluminescence at 430 nm from the commercial chip further excites Sn-doped CsBr and gives rise to its characteristic emission identical to the PL. In the inset, a zoom view of the emission region belonging to the Sn-doped CsBr is shown. A Commission Internationale de l'Éclairage color coordinate is provided in Figure S15 in the Supporting Information to show the contribution of orange emission. The

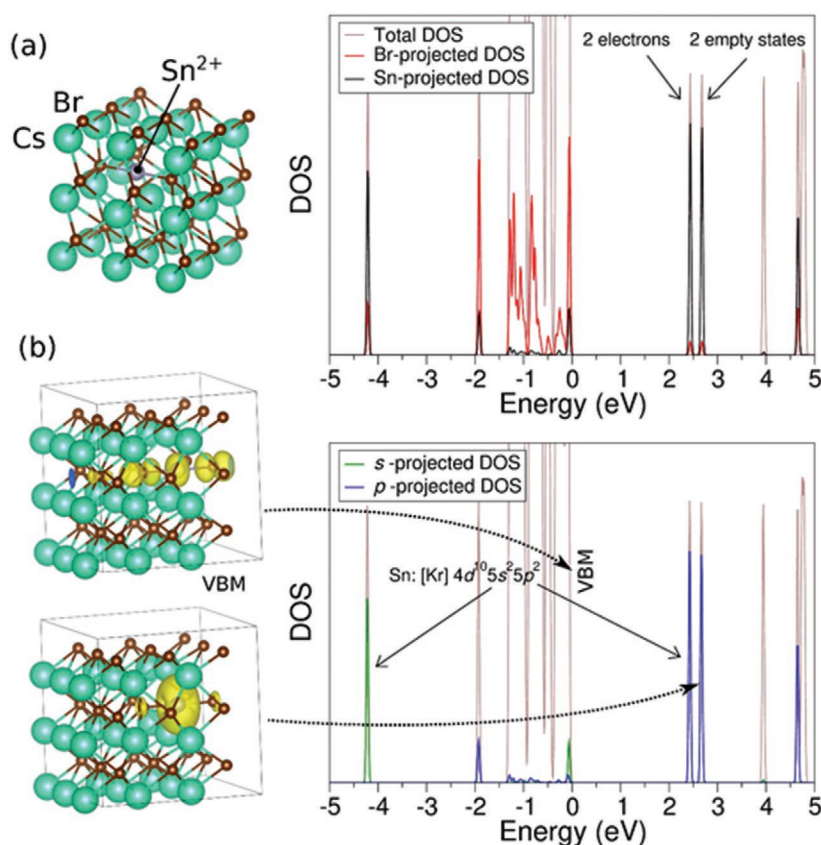


Figure 4. a) Supercell of cubic CsBr with an interstitial Sn dopant centered in an orthorhombic void with two Br atoms aside. Total densities of states (DOS) with indication of the Sn contribution. b) Total DOS with indication of the *sp*-contributions of the dopant. Isosurfaces of constant charge density for a self-trapped exciton composed of a hole at the valence band maximum (VBM, upper panel) and the promoted electron occupying one of the available impurity states (lower panel).

broad spectral characteristics of the PL signals and the moderate to good PLQY values obtained from the Sn-doped CsBr material, together with the outstanding (photo-)chemical stability suggests that this material could be a good candidate to be exploited as a down-converter toward the developing of white light-emitting devices. Additionally, our strategy opens up a new paradigm for developing a new family of materials, through a systematic compositional screening, suitable for optoelectronic applications.

In order to qualitatively quantify the scope of the continuous flow reaction methodology in particular study, a parallel reaction was carried out manually by the dropwise addition of 0.5 mL of precursor solution to 1 mL of antisolvent (toluene) under continuous stirring conditions. The product obtained by this manual procedure is also characterized as Sn(II)-doped CsBr from the XRD. In comparison to the flow reactor method, the material obtained from the manual reaction procedure has less optical performance, and poor particle shape definition

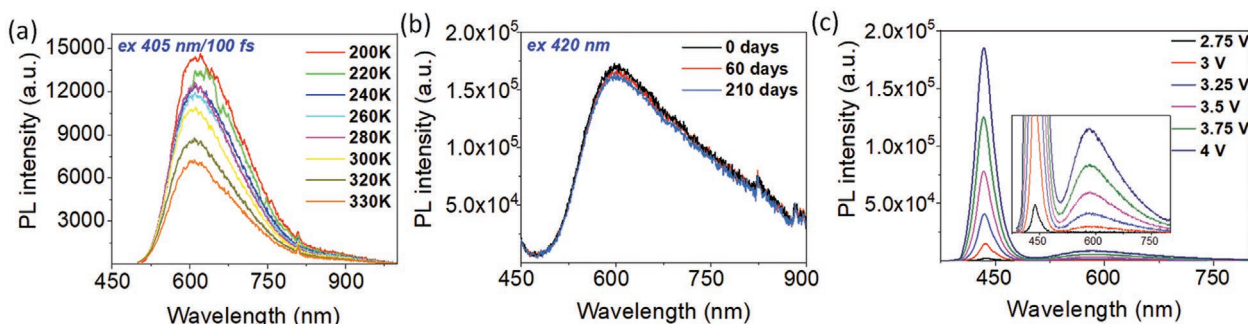


Figure 5. a) PL spectra of Sn(II)-doped CsBr measured under 405 nm excitation by 200 fs laser pulses at different temperatures from 200 to 330 K. b) PL spectra of Sn(II)-doped CsBr measured at different times after synthesis. c) Down-converted PL spectra of Sn(II)-doped CsBr excited by 430 nm electroluminescence of a commercial blue chip made of (InGa)N.

and size distribution. A comparison to the optical emission and the SEM images is provided in Figures S16 and S17 in the Supporting Information. In Figure S17 in the Supporting Information, SEM image of undoped CsBr shows an identical shape and size distribution to the Sn(II)-doped CsBr prepared using flow-reactor. This is due to the improved mixing achieved under continuous-flow conditions, which is particularly relevant for fast reactions. This procedure allows to generate advanced nanostructured materials under steady state. Hence, the flow reactor is much more effective toward i) production of mass-scale continuously, ii) formation of material at short-time scale leads to the formation of equal-shaped nanocrystals, and iii) enhanced optical performance of crystalline products.

3. Conclusion

In summary, we have prepared a new material belonging to the lead-free cesium metal halide family, which is prepared using a continuous flow reactor. By altering the precursor flow rates, which is indeed the ratio of precursor concentration to the anti-solvent, we obtained a tin-rich phase Cs₄SnBr₆ at low precursor concentration. With the increase of precursor concentration by increasing the precursor flow rate, a lesser amount of tin was precipitated along with the CsBr, which led to the formation of the new material, Sn(II)-doped CsBr. This material possesses a broadband STE emission at the orange-red region, which emits PL with a PLQY up to 21.5% if excited in the 390–435 nm diapason. This material has a visible-light absorbance near 2.8 eV due to mid-gap localized states contributed by the dopant which were further confirmed by quantum mechanical calculations based on the density functional theory. It is worth to highlight, that Sn(II)-doped CsBr is highly stable in the open air, and its emission was found stable over 7 months. The present state-of-art of this finding leads us to the proof of concept for the lead-free cesium metal halides as a light-emitting (down-shift wavelength converter) material. The broadband emission from this material could be implemented for white-LED applications.

4. Experimental Section

Preparation of Precursor Solution: 212 mg (1 mmol) of CsBr and 278 mg (1 mmol) of SnBr₂ were added to 25 mL of DMF and stirred for 30 min to obtain a clear solution. Finally, this solution was used as the precursor solution.

Flow-Reactor Synthesis: Two different flow pumps were used and connected with a “T” junction. Herein, toluene was used as the anti-solvent. Toluene was pumped using one of the pumps connected parallel to the product-side capillary, and precursor solution was pumped using another

pump, which was connected perpendicular to the product-side capillary. The flow-rate of toluene was fixed to 1 mL min⁻¹ and the flow-rate of precursor solution was changed from 0.1 to 0.5 mL min⁻¹. Details of flow-rate and obtained different products are summarized in **Table 1**.

Preparation of Cs₂SnBr₆: 212 mg of CsBr and 278 mg of SnBr₂ were added in 2 mL of acetone and stirred for overnight at 50 °C. A brown colored powder was formed.

PMMA-Coating for XRD Characterization of As-Prepared Cs₄SnBr₆: At 0.1/1 P/A, the green emitting material was collected under N₂ atmosphere after 10 min of flow (total volume = 11 mL), finally the obtained material was kept inside the glove box in an air-tightened glass vial and settle the product + solution. After 45–60 min, the supernatant was decanted and the precipitate was collected. In another vial, 100 mg of PMMA and 1 mL of chlorobenzene was loaded together, heated to 60 °C for 30 min, poured to the collected Cs₄SnBr₆ product once a clear solution of PMMA obtained. It was stirred for another 30 min. Finally, the product embedded in PMMA+chlorobenzene solution spin-coated on top of a glass slide, dried to ambient condition. Noteworthy, all these procedures were carried out inside the glove box. This film was also found to be stable over 7 months.

Characterizations: UV–vis absorption spectra were collected from Perkin Elmer UV–vis–NIR spectrophotometer, the powder samples were measured using an integrated sphere unit. PL and PLQY measurements were carried out using an integrating sphere unit (Hamamatsu). The samples were measured in a powder form. The Sn(II)-doped CsBr samples were excited at 420 nm and Cs₄SnBr₆ samples were excited at 340 nm. PLE was measured using fluorimeter Horiba. XRD measurements were performed using D8 Advance, Bruker-AXS X-ray diffractometer. The powder samples were dispersed in toluene and dropcasted onto a glass slide for the XRD measurement (Cu K α , wavelength λ = 1.5406 Å). Microscopic images were captured using SEM (JSM-7001F, JEOL) and TEM (JEOL, JEM-2100). EDS analysis was carried out from TEM measurements. TGA was measured using TGA/SDTA851e/LF/1600 instrument and the temperature range was from ambient temperature to 1000 °C. The temperature increment rate was 5 °C min⁻¹. RAMAN spectroscopic measurement for the powder samples was carried out using RAMAN WITEC APYRON instrument, the laser wavelength for all the samples were 785 nm.

XPS measurements were performed using a Kratos AXIS Ultra DLD instrument. The chamber pressure during the measurements was 5 × 10⁻⁹ Torr. Wide energy range survey scans were collected at a pass energy of 80 eV in hybrid slot lens mode and a step size of 0.5 eV, for 20 min. The charge neutralizer filament was used to prevent the sample charging over the irradiated area. The X-ray source was a monochromated Al K α emission, run at 10 mA and 12 kV (120 W). The energy range for each “pass energy” (resolution) was calibrated using the Kratos Cu 2p 3/2, Ag 3d 5/2, and Au 4f 7/2 three-point calibration method. The transmission function was calibrated using a clean gold sample method for all lens modes and the Kratos transmission generator software within Vision II. The data were processed with CASAXPS (Version 2.3.17). The high-resolution data were charge corrected to the reference C 1s signal at 285 eV.

For low temperature PL and time-resolved PL (TRPL) measurements, the samples were held in the coldfinger of a closed-cycle He cryostat (ARS DE-202), which can be cooled down to 15 K. PL was excited by means of 200 fs pulsed Ti:sapphire passive mode-locked laser (Coherent

Table 1. Details of flow-rates and obtained materials.

Entry	Precursor flow-rate [mL min ⁻¹]	Antisolvent flow-rate [mL min ⁻¹]	Obtained material	Emission under 365 nm UV lamp	Stability
1	0.1	1	Cs ₄ SnBr ₆	Green	Few hours
2	0.2	1	Cs ₄ SnBr ₆	Green	Few hours
3	0.3	1	Sn(II)-doped CsBr	Orange	Over 7 months
4	0.4	1	Sn(II)-doped CsBr	Orange	Over 7 months
5	0.5	1	Sn(II)-doped CsBr	Orange	Over 7 months

Mira 900D, 76 MHz repetition rate) operating at a wavelength of 810 nm and doubled to 405 nm by using a BBO crystal. The PL signal was dispersed by a double 0.3 m focal length grating spectrograph (Acton SP-300i from Princeton Instruments) and detected with a cooled Si charge-coupled device camera (Newton EMCCD from ANDOR) for PL spectral measurements and with a silicon single photon avalanche photodiode (micro photon device) connected to a time-correlated single-photon counting electronic board (TCC900 from Edinburgh Instruments) for TRPL measurements.

Supporting Information

Supporting Information is available from the Wiley Online Library or from the author.

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Conflict of Interest

The authors declare no conflict of interest.

Data Availability Statement

Research data are not shared.

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continuous flow reactor synthesis, CsBr, interstitial doping, long-term ambient stability, photoluminescence quantum yield, self-trapped excitonic emission

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- [1] J. Jeong, M. Kim, J. Seo, H. Lu, P. Ahlawat, A. Mishra, Y. Yang, M. A. Hope, F. T. Eickemeyer, M. Kim, Y. J. Yoon, I. W. Choi, B. P. Darwich, S. J. Choi, Y. Jo, J. H. Lee, B. Walker, S. M. Zakeeruddin, L. Emsley, U. Rothlisberger, A. Hagfeldt, D. S. Kim, M. Grätzel, J. Y. Kim, *Nature* **2021**, 592, 381.
- [2] W. Xu, Q. Hu, S. Bai, C. Bao, Y. Miao, Z. Yuan, T. Borzda, A. J. Barker, E. Tyukalova, Z. Hu, M. Kawecki, H. Wang, Z. Yan, X. Liu, X. Shi, K. Uvdal, M. Fahlman, W. Zhang, M. Duchamp, J.-M. Liu, A. Petrozza, J. Wang, L.-M. Liu, W. Huang, F. Gao, *Nat. Photonics* **2019**, 13, 418.
- [3] C. Chen, D. Li, Y. Wu, C. Chen, Z.-G. Zhu, W. Y. Shih, W.-H. Shih, *Nanotechnology* **2020**, 31, 225602.

- [4] C. Chen, Y. Wu, Z.-G. Zhu, W. Y. Shih, W.-H. Shih, *J. Mater. Res.* **2021**, 36, 1835.
- [5] J. Lin, Y. Lu, X. Li, F. Huang, C. Yang, M. Liu, N. Jiang, D. Chen, *ACS Energy Lett.* **2021**, 6, 519.
- [6] S. Yuan, D. Chen, X. Li, J. Zhong, X. Xu, *ACS Appl. Mater. Interfaces* **2018**, 10, 18918.
- [7] J. Li, H.-L. Cao, W.-B. Jiao, Q. Wang, M. Wei, I. Cantone, J. Lü, A. Abate, *Nat. Commun.* **2020**, 11, 310.
- [8] M. Lyu, J.-H. Yun, P. Chen, M. Hao, L. Wang, *Adv. Energy Mater.* **2017**, 7, 1602512.
- [9] S. M. Jain, T. Edvinsson, J. R. Durrant, *Commun. Chem.* **2019**, 2, 91.
- [10] P. V. Kamat, J. Bisquert, J. Buriak, *ACS Energy Lett.* **2017**, 2, 904.
- [11] B. Saparov, F. Hong, J.-P. Sun, H.-S. Duan, W. Meng, S. Cameron, I. G. Hill, Y. Yan, D. B. Mitzi, *Chem. Mater.* **2015**, 27, 5622.
- [12] A. H. Slavney, T. Hu, A. M. Lindenberg, H. I. Karunadasa, *J. Am. Chem. Soc.* **2016**, 138, 2138.
- [13] T.-B. Song, T. Yokoyama, S. Aramaki, M. G. Kanatzidis, *ACS Energy Lett.* **2017**, 2, 897.
- [14] H. Chen, J. M. Pina, F. Yuan, A. Johnston, D. Ma, B. Chen, Z. Li, A. Dumont, X. Li, Y. Liu, S. Hoogland, Z. Zajac, Z. Lu, E. H. Sargent, *J. Phys. Chem. Lett.* **2020**, 11, 4326.
- [15] J. Luo, X. Wang, S. Li, J. Liu, Y. Guo, G. Niu, L. Yao, Y. Fu, L. Gao, Q. Dong, C. Zhao, M. Leng, F. Ma, W. Liang, L. Wang, S. Jin, J. Han, L. Zhang, J. Etheridge, J. Wang, Y. Yan, E. H. Sargent, J. Tang, *Nature* **2018**, 563, 541.
- [16] J. Li, J. Duan, X. Yang, Y. Duan, P. Yang, Q. Tang, *Nano Energy* **2021**, 80, 105526.
- [17] Z. Zhang, R. Zhao, S. Teng, K. Huang, L. Zhang, D. Wang, W. Yang, R. Xie, N. Pradhan, *Small* **2020**, 16, 2004272.
- [18] E. López-Fraguas, S. Masi, I. Mora-Seró, *ACS Appl. Energy Mater.* **2019**, 2, 8381.
- [19] S. Masi, A. F. Gualdrón-Reyes, I. Mora-Seró, *ACS Energy Lett.* **2020**, 5, 1974.
- [20] Z. Wang, J. Zhang, W. Guo, W. Xiang, A. Hagfeldt, *Matter* **2021**, 4, 528.
- [21] J. Cao, F. Yan, *Energy Environ. Sci.* **2021**, 14, 1286.
- [22] X. Li, X. Gao, X. Zhang, X. Shen, M. Lu, J. Wu, Z. Shi, V. L. Colvin, J. Hu, X. Bai, W. W. Yu, Y. Zhang, *Adv. Sci.* **2021**, 8, 2003334.
- [23] M. Li, Z. Xia, *Chem. Soc. Rev.* **2021**, 50, 2626.
- [24] K. M. McCall, V. Morad, B. M. Benin, M. V. Kovalenko, *ACS Mater. Lett.* **2020**, 2, 1218.
- [25] A. E. Maughan, A. M. Ganose, D. O. Scanlon, J. R. Neilson, *Chem. Mater.* **2019**, 31, 1184.
- [26] T. C. Jellicoe, J. M. Richter, H. F. J. Glass, M. Tabachnyk, R. Brady, S. E. Dutton, A. Rao, R. H. Friend, D. Credgington, N. C. Greenham, M. L. Böhm, *J. Am. Chem. Soc.* **2016**, 138, 2941.
- [27] B. M. Benin, D. N. Dirin, V. Morad, M. Wörle, S. Yakunin, G. Rainò, O. Nazarenko, M. Fischer, I. Infante, M. V. Kovalenko, *Angew. Chem., Int. Ed.* **2018**, 57, 11329.
- [28] R. Chiara, Y. O. Ciftci, V. I. E. Queloz, M. K. Nazeeruddin, G. Grancini, L. Malavasi, *J. Phys. Chem. Lett.* **2020**, 11, 618.
- [29] X. Zhang, H. Wang, S. Wang, Y. Hu, X. Liu, Z. Shi, V. L. Colvin, S. Wang, W. W. Yu, Y. Zhang, *Inorg. Chem.* **2020**, 59, 533.
- [30] L. Tan, W. Wang, Q. Li, Z. Luo, C. Zou, M. Tang, L. Zhang, J. He, Z. Quan, *Chem. Commun.* **2020**, 56, 387.
- [31] V. Santhana, D. C. Greenidge, D. Thangaraju, R. Marnadu, T. Alshahrani, M. Shkir, *Mater. Lett.* **2020**, 280, 128562.
- [32] S. Li, J. Luo, J. Liu, J. Tang, *J. Phys. Chem. Lett.* **2019**, 10, 1999.
- [33] A. Veronese, M. Patrini, D. Bajoni, C. Ciarrocchi, P. Quadrelli, L. Malavasi, *Front. Chem.* **2020**, 8, 35.
- [34] I. Lignos, S. Stavrakis, G. Nedelcu, L. Protesescu, A. J. deMello, M. V. Kovalenko, *Nano Lett.* **2016**, 16, 1869.

- [35] R. M. Maceiczky, K. Dümbgen, I. Lignos, L. Protesescu, M. V. Kovalenko, A. J. deMello, *Chem. Mater.* **2017**, *29*, 8433.
- [36] I. Lignos, L. Protesescu, D. B. Emiroglu, R. Maceiczky, S. Schneider, M. V. Kovalenko, A. J. deMello, *Nano Lett.* **2018**, *18*, 1246.
- [37] X. Liang, R. W. Baker, K. Wu, W. Deng, D. Ferdani, P. S. Kubiak, F. Marken, L. Torrente-Murciano, P. J. Cameron, *React. Chem. Eng.* **2018**, *3*, 640.
- [38] O. Okafor, K. Robertson, R. Goodridge, V. Sans, *React. Chem. Eng.* **2019**, *4*, 1682.
- [39] O. Okafor, A. Weilhard, J. A. Fernandes, E. Karjalainen, R. Goodridge, V. Sans, *React. Chem. Eng.* **2017**, *2*, 129.
- [40] V. Sans, S. Glatzel, F. J. Douglas, D. A. Maclaren, A. Lapkin, L. Cronin, *Chem. Sci.* **2014**, *5*, 1153.
- [41] S. Yakunin, B. M. Benin, Y. Shynkarenko, O. Nazarenko, M. I. Bodnarchuk, D. N. Dirin, C. Hofer, S. Cattaneo, M. V. Kovalenko, *Nat. Mater.* **2019**, *18*, 846.
- [42] A. Kaltzoglou, M. Antoniadou, A. G. Kontos, C. C. Stoumpos, D. Perganti, E. Siranidi, V. Raptis, K. Trohidou, V. Psycharis, M. G. Kanatzidis, P. Falaras, *J. Phys. Chem. C* **2016**, *120*, 11777.
- [43] T.-B. Song, T. Yokoyama, C. C. Stoumpos, J. Logsdon, D. H. Cao, M. R. Wasielewski, S. Aramaki, M. G. Kanatzidis, *J. Am. Chem. Soc.* **2017**, *139*, 836.
- [44] P. Giannozzi, O. Andreussi, T. Brumme, O. Bunau, M. Buongiorno Nardelli, M. Calandra, R. Car, C. Cavazzoni, D. Ceresoli, M. Cococcioni, N. Colonna, I. Carnimeo, A. Dal Corso, S. de Gironcoli, P. Delugas, R. A. DiStasio, A. Ferretti, A. Floris, G. Fratesi, G. Fugallo, R. Gebauer, U. Gerstmann, F. Giustino, T. Gorni, J. Jia, M. Kawamura, H. Y. Ko, A. Kokalj, E. Küçükbenli, M. Lazzeri, M. Marsili, N. Marzari, F. Mauri, N. L. Nguyen, H. V. Nguyen, A. Otero-de-la-Roza, L. Paulatto, S. Poncé, D. Rocca, R. Sabatini, B. Santra, M. Schlipf, A. P. Seitsonen, A. Smogunov, I. Timrov, T. Thonhauser, P. Umari, N. Vast, X. Wu, S. Baroni, *J. Phys.: Condens. Matter* **2017**, *29*, 465901.
- [45] K. Teegarden, G. Baldini, *Phys. Rev.* **1967**, *155*, 896.